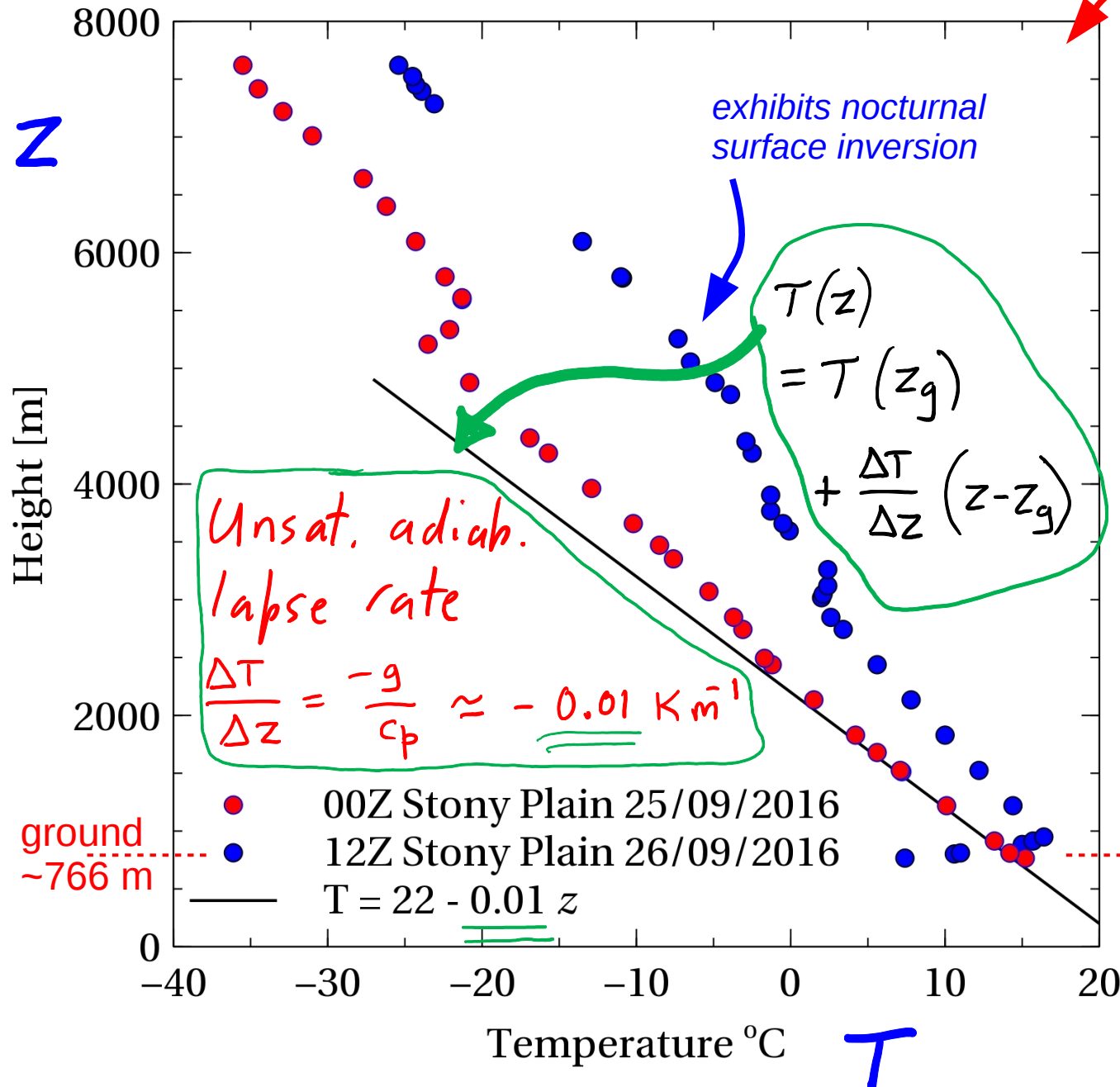


Aside: an overbar means an average of some kind, e.g.  $\bar{T}$

Demonstration of the practical relevance of the adiabatic lapse rate



### Afternoon sounding

Shallow surface layer has lapse rate stronger than the adiabatic rate

This is due to surface heating

Above the surface layer, lapse rate is adiabatic to about 2500 m AGL

ground level  
~766 m

say " $z_g$ "

$$\Delta q = 0 = c_p \Delta T - \frac{\Delta P}{\rho}$$

$$P = \rho R_d T$$

$$T_1, P_1 \xrightarrow{\text{adiabatic process}} T, P$$

$$\frac{1}{\rho} = \frac{R_d T}{P}$$

$$c_p \Delta T = R_d T \frac{\Delta P}{P}$$

$$\frac{\Delta T}{T} = \frac{R_d}{c_p} \frac{\Delta P}{P}$$

fractional change

$$\Delta \ln T = \frac{R_d}{c_p} \Delta \ln P$$

$$\frac{\Delta \ln T}{\Delta \ln P} = \frac{R_d}{c_p}$$

Upon integration from state  $(P_1, T_1)$  to  $(P, T)$  one obtains

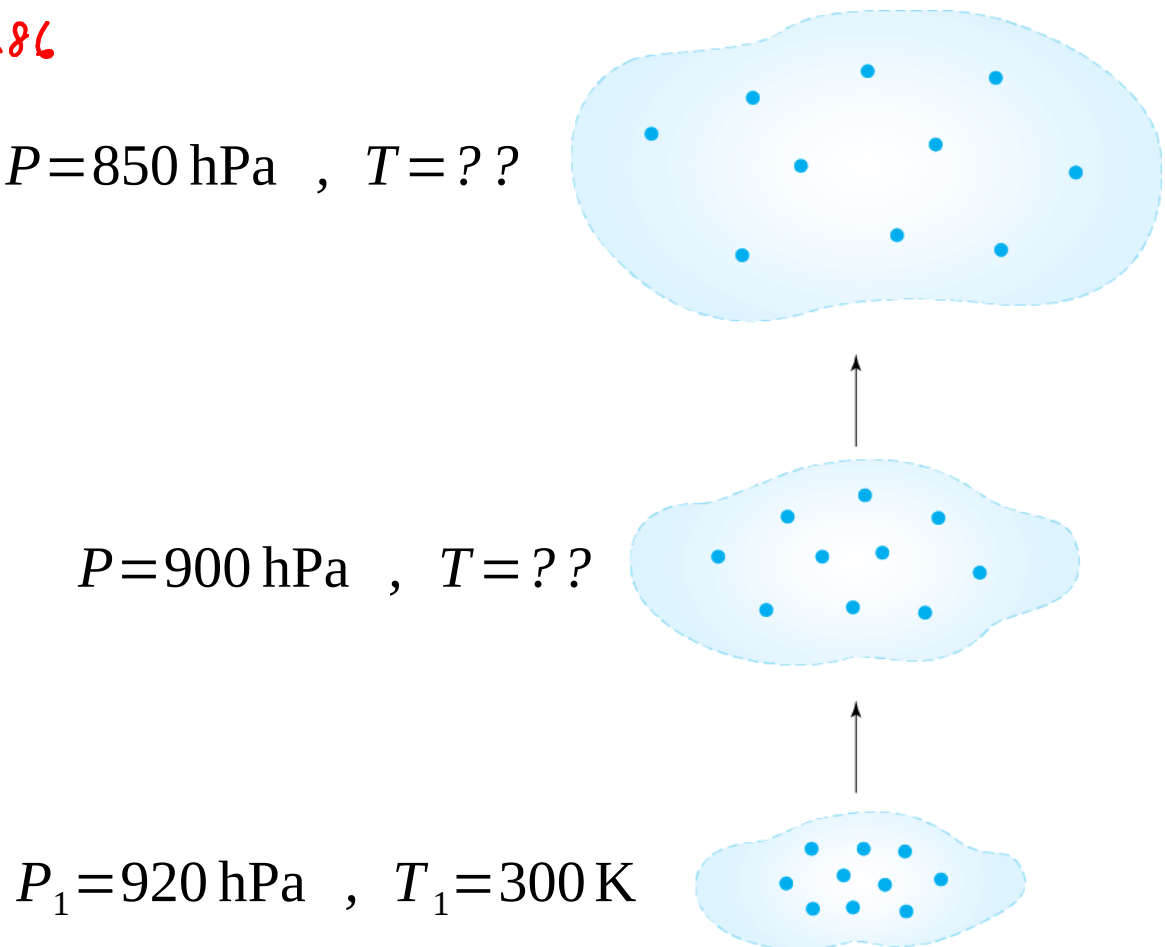
$$\frac{T}{T_1} = \left[ \frac{P}{P_1} \right]^{\frac{R_d}{c_p}}$$

where  $\frac{R_d}{c_p} = \frac{2}{7} = 0.286$

$$T/T_1 = (P/P_1)^{0.286}$$

Poisson's equation is useful if we want to compute the height variation of temperature based on ascent/descent across specified pressure levels (rather than specified elevations). Suppose the surface parcel rises adiabatically: compute its temperatures at 900 hPa and at 850 hPa...

$$T = T_1 \left( \frac{P}{P_1} \right)^{0.286} = 300 \left( \frac{850}{920} \right)^{0.286} \quad P=850 \text{ hPa} , T=? ?$$
$$= 293.29$$



Quantity of sensible heat per unit volume\*\*, say “ $h$ ” [ $\text{J m}^{-3}$ ], is given by:

$$h = \rho c_p T$$

$\rho$  ← air density  
 $c_p$  ← specific heat capacity of air  
 $\sim 1000 \text{ J kg}^{-1} \text{ K}^{-1}$   
 $T$  ← temperature, K

Let  $\rho_v [\text{kg/m}^3]$   
 be amount of  
 vapour per unit vol.

\*\* Alternative name: “volumetric (sensible) heat content”. This is the appropriate measure in the context of heat transfer occurring at constant pressure, recalling

$$\Delta q = c_p \Delta T - \alpha \Delta P$$

Quantity of *latent* heat per unit volume, say “ $h_L$ ” [ $\text{J m}^{-3}$ ], is given by:

$$h_L = \rho_v L_v = \text{absolute humidity of the air } [\text{kg m}^{-3}] \times \text{latent heat of vapourization } \sim 2.5 \times 10^6 \text{ J kg}^{-1}$$

molecular level processes

## A. Conduction\*\* / Diffusion

- occurs in solids, liquids & gases
- spatial difference in temperature (mean kinetic energy of material particles) implies energetically vibrating matter in contact with less energetically vibrating material
- vibration energy is transferred
- football crowd analogy

- Fourier's law of conduction  $Q_x = -k \frac{dT}{dx}$

Flux of sensible heat ( $Q_x$ ,  $W m^{-2}$ ) along direction  $x$  equals (minus) the conductivity ( $k$ ) times the spatial gradient of temperature in the  $x$  direction.

Vector generalisation:

$$\vec{Q} = (Q_x, Q_y, Q_z) = -k \left( \frac{dT}{dx}, \frac{dT}{dy}, \frac{dT}{dz} \right)$$

What are the units of  $k$ ?

$$[k] = \frac{[Q]}{[dT/dx]} = \frac{W m^{-2}}{K m^{-1}} = W m^{-1} K^{-1}$$

\*\*Similar mechanisms and laws for mass "diffusion" and viscous momentum transfer

## B. Convection

Transport by bulk flow

wind

(Fick's law)

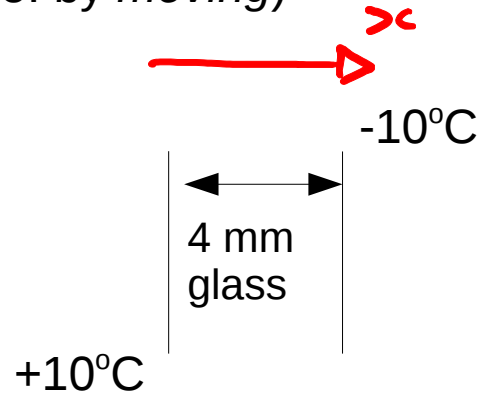
(Newton's law of viscosity)

\*Sec 4.4 is titled "Heat Transfer", but we'll cover the broader topic of "Natural mechanisms for heat, mass and momentum transfer"

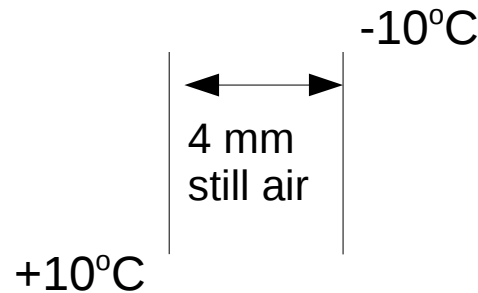
**C. Radiation (subject of Chapter 5)** transports solar energy across the vacuum of space and powers the earth; we also deal with radiation from (and within) the earth-atmosphere system itself (terrestrial radiation). We later learn how greenhouse gases operate as selective absorbers of terrestrial (longwave) radiant energy

- by definition, poor conductors are good insulators – still air or fresh snow
- in the atmosphere, conduction is only important very close to surfaces (e.g. ground, or leaf surfaces) within the "laminar boundary layer". Air transports heat more effectively by convection (i.e. by *moving*)

Example:



$$|Q| \approx k \frac{\Delta T}{\Delta x} \approx 1 \frac{20}{.004} = 5000 \text{ W m}^{-2}$$



$$|Q| \approx .025 \times \frac{20}{.004} = 125 \text{ W m}^{-2}$$

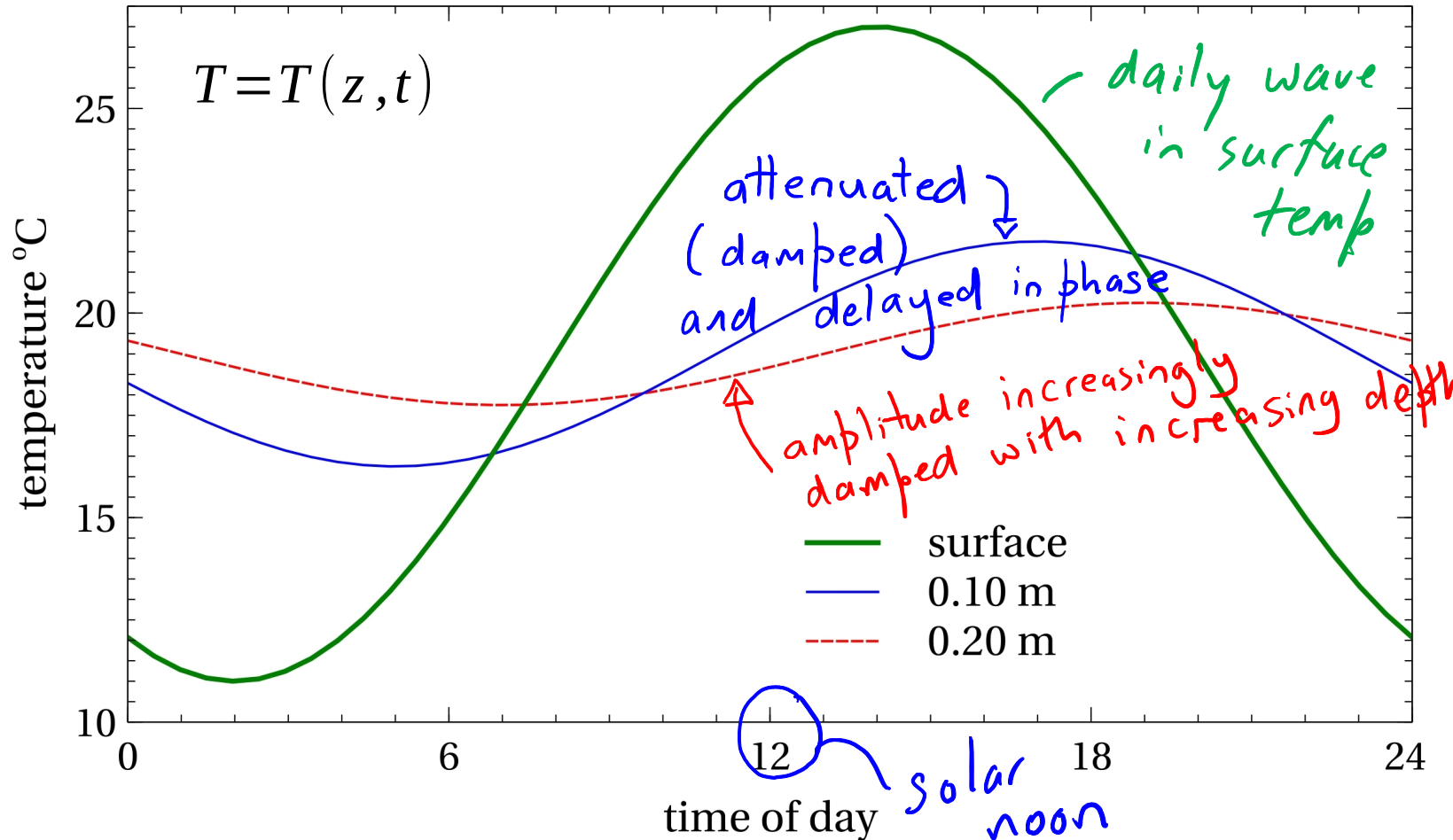
**TABLE 4.2 | Conductivities of various substances.**

Substance	Conductivity ( $\text{W}\cdot\text{m}^{-1}\cdot\text{K}^{-1}$ )
Still Air ( $10^\circ\text{C}$ )	0.025
Water ( $10^\circ\text{C}$ )	0.62
Ice ( $0^\circ\text{C}$ )	2.24
Fresh Snow	0.08
Old Snow	0.42
Dry Sand	0.15–0.25
Moist Sand	0.25–2.0
Granite	2.9
Limestone	1.3
Light Wood	0.09
Dense Wood	0.19
Stainless Steel	16
Aluminum	200
Glass	~1

Generally higher for solids, due to their higher density (atoms or molecules in close proximity)

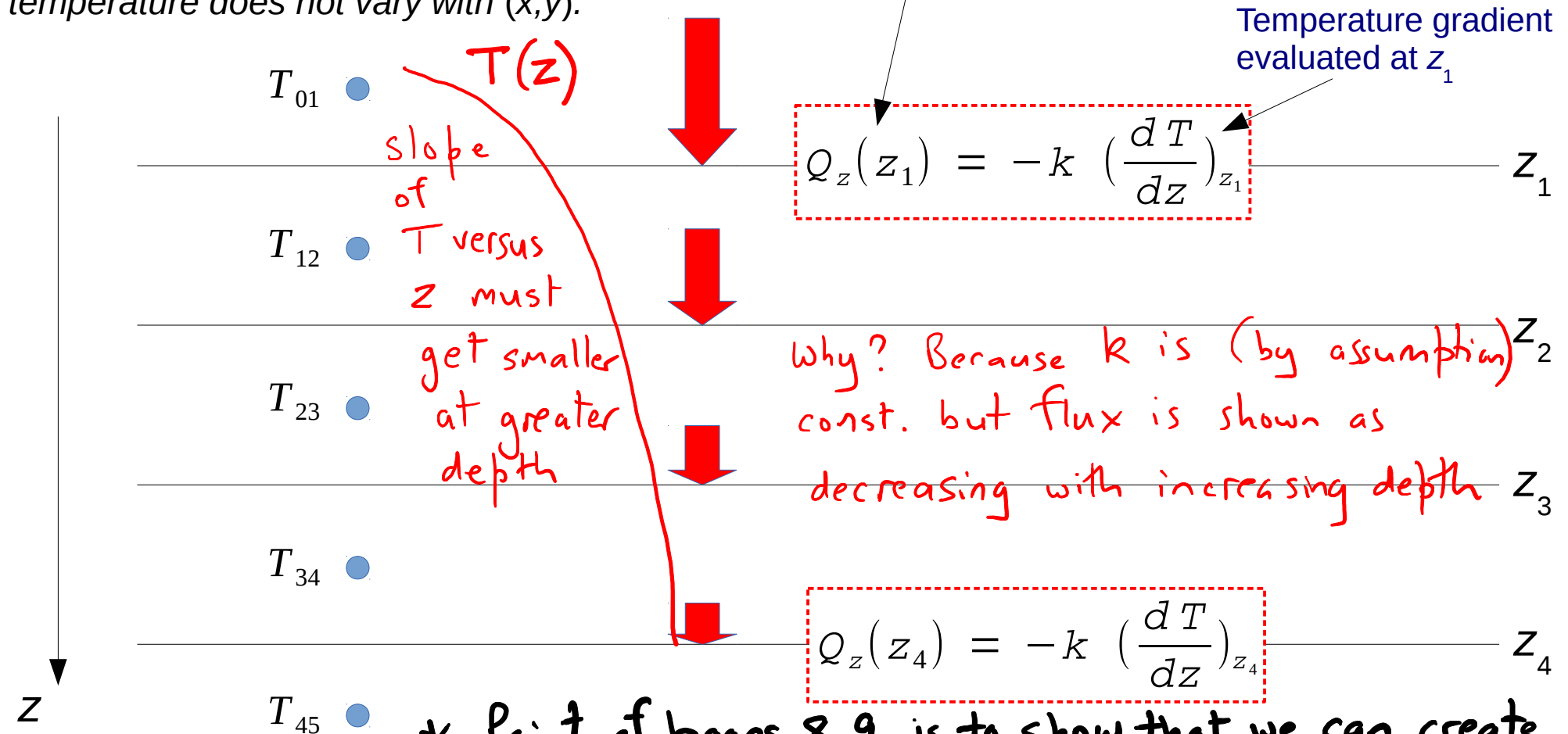
- conduction is the main mechanism for heat transfer within ground (percolation of rainwater – a convective process – can also transfer heat)
- daily (diurnal) and annual "temperature waves" are attenuated with increasing depth (  $z$  )
- ~ 10 cm participates diurnally; ~ few m annually
- soil conductivity a strong function of soil moisture content

Penetration of the daily temperature wave – increasing attenuation (i.e. amplitude reduction) and phase lag (i.e. time delay of the temperature peak) with increasing depth.



- consider layers of soil, assuming soil conductivity  $k$  is uniform (constant). The soil heat flux  $Q_z$  varies with depth  $z$ :

*Implicitly we're assuming horizontal heat transport is negligible. This amounts to an assumption of symmetry, i.e. that soil temperature does not vary with (x,y).*



\* Point of pages 8,9 is to show that we can create a quantitative model of the "meteorology" of the soil (i.e. the environmental state).



In time  $\Delta t$  the volume  $A \Delta z$  of the layer whose temperature is  $T_{23}$  gains an amount of heat  $Q_z(z_2) A \Delta t$  [J] conducted across  $z_2$ , but loses heat  $Q_z(z_3) A \Delta t$  [J] conducted across  $z_3$

Soil density and heat capacity  $\rho_s$  and  $c_s$

$\text{kg m}^{-3}$   $\text{J kg}^{-1} \text{K}^{-1}$

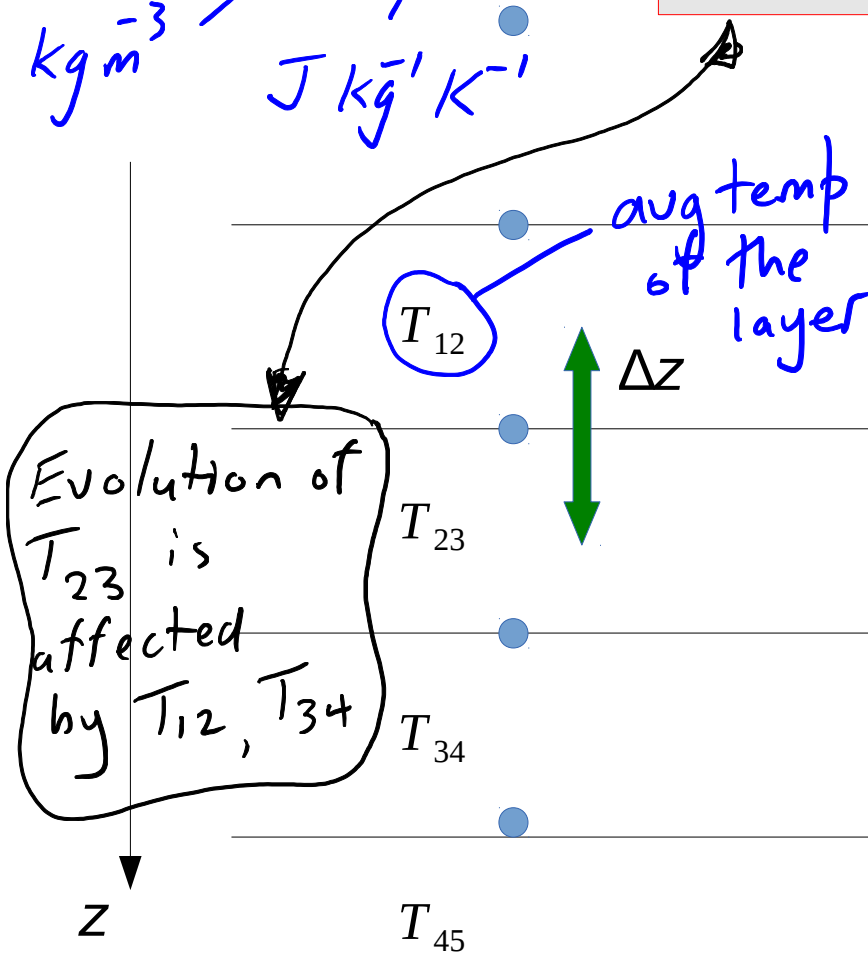
$$\rho_s c_s \underbrace{A \Delta z}_{\text{volume}} \Delta T_{23} = A \underbrace{\Delta t}_{\text{time step}} [Q_z(z_2) - Q_z(z_3)]$$

We shan't take any further *etc.* this model, suffice to say that one could. Already we see the inter-connectivity of the

$$Q_z(z_2) = -k \left( \frac{dT}{dz} \right)_{z_2} = -k \frac{T_{23} - T_{12}}{\Delta z}$$

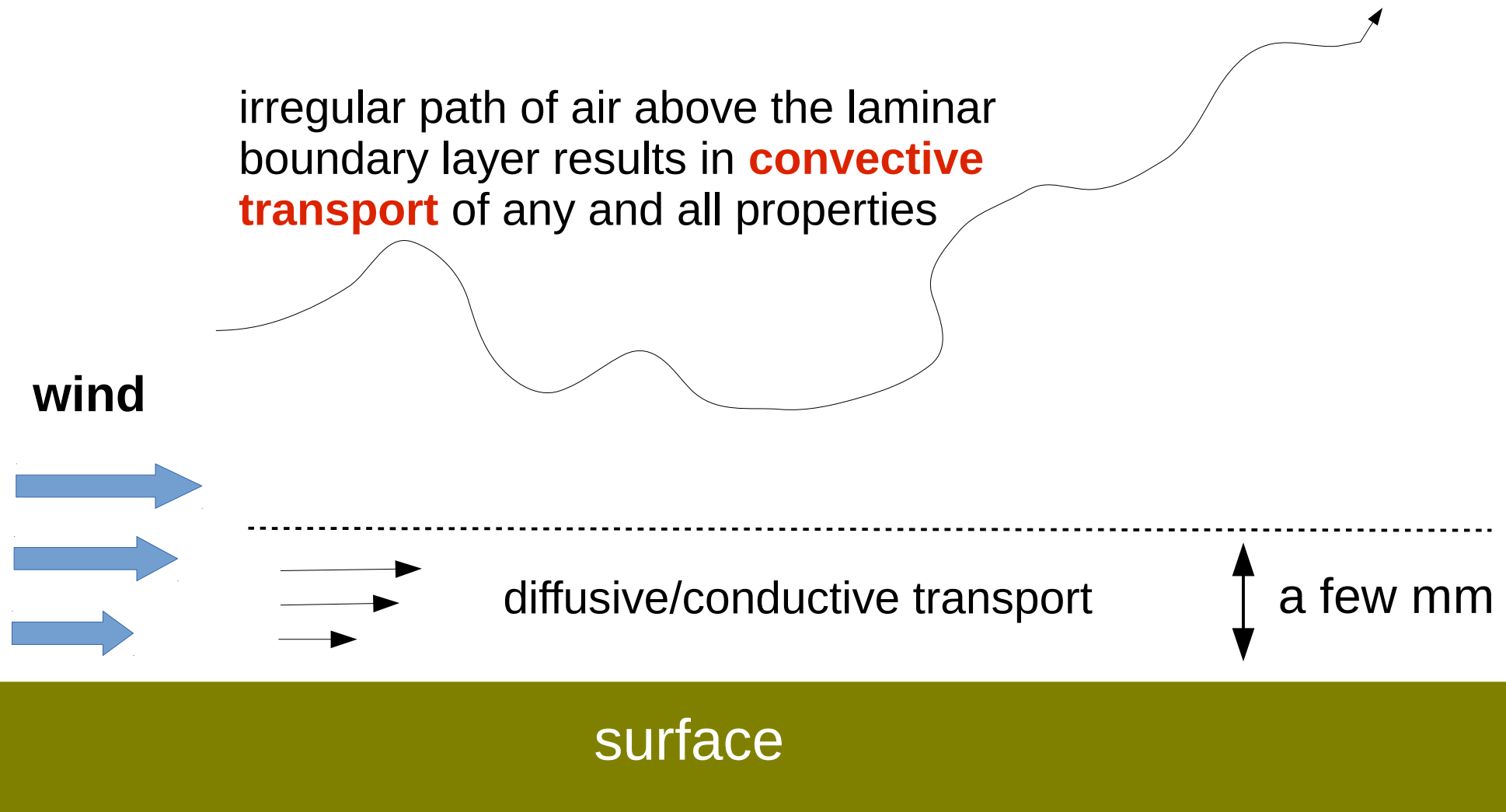
Soil environment

$$Q_z(z_4) = -k \left( \frac{dT}{dz} \right)_{z_4}$$



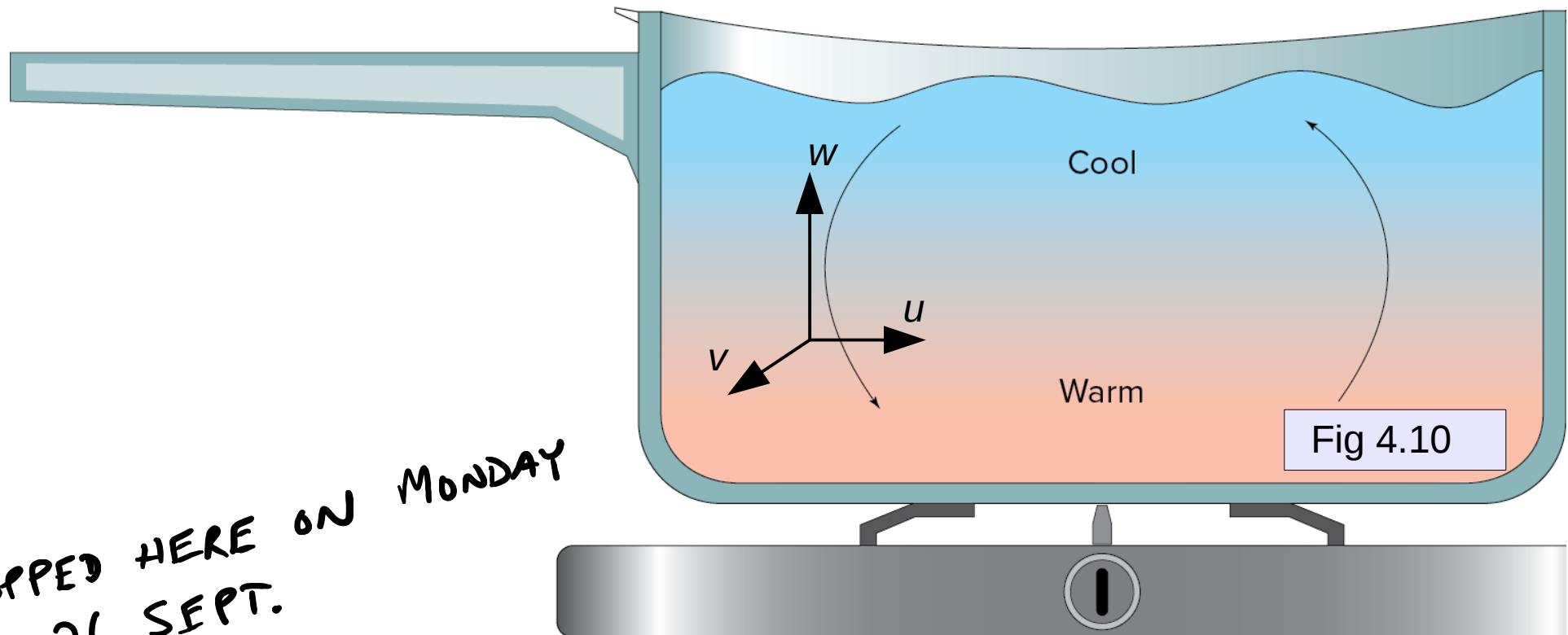
- at each interface, we can evaluate the needed heat flux using Fourier's law

When a fluid or gas flows over a solid surface, there is a thin layer adjacent to the surface, known as the laminar boundary layer, in which the flow is *parallel* to the surface, meaning there can be no convective transport to or from that surface – in that layer, only radiation & conduction can transport heat, and only diffusion can transport mass (e.g. water vapour)



*wind*

- Transport of heat by "mass movement" of a fluid, i.e. organized (bulk) flow as opposed to the jitter of individual atoms, is named "convection"
- In the broadest context convection may transport *any* property in *any* direction
- Commonly the word is used to refer (more narrowly) to vertical transport of heat
- Must interpret the words "convection, convective" according to the context



STOPPED HERE ON MONDAY  
26 SEPT.

- always wind velocity x volumetric concentration <sup>amount per cubic metre</sup> of the transported property
- e.g. let  $w$  be the vertical wind speed and  $h$  (as before) be the volumetric sensible heat content, then the rate of vertical convective transport of sensible heat is simply:

$$\underline{Q_z} = w h = \rho c_p w T \quad [\text{J s}^{-1} \text{m}^{-2} = \text{W m}^{-2}]$$

- this quantifies the vertical convection of energy in one form, i.e. sensible heat energy. To use all the adjectives, this is: the instantaneous local vertical convective flux density of sensible heat
- it's an important quantity that can be, and often is, measured (more in a later chapter). What is important here is the principle: product of wind velocity (in a given direction) times the amount of “stuff” per unit volume of air gives the rate at which that “stuff” is being transported (its “flux” in the given direction)

Q: why is the sensible heat flux density evaluated using  $c_p$  rather than  $c_v$ ?

because heat exchange is occurring at const. pressure

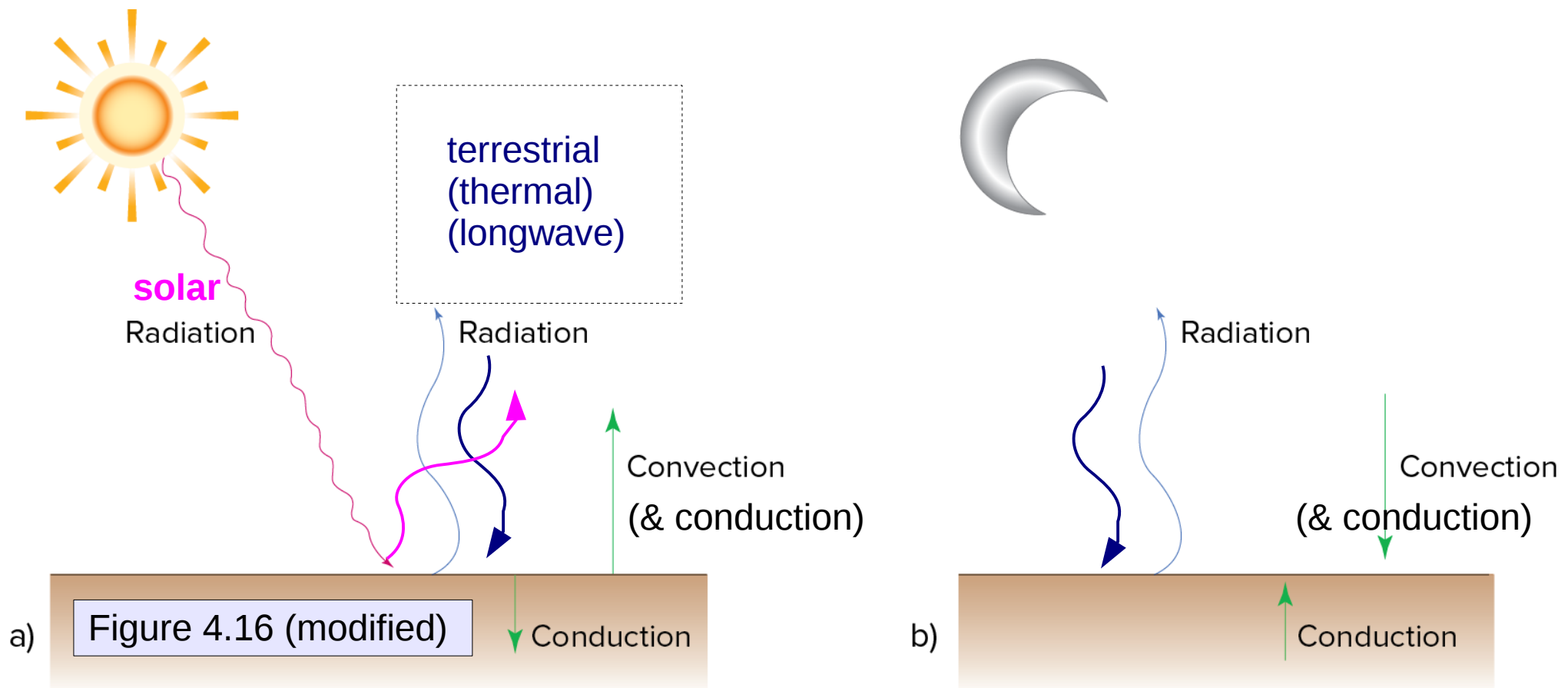
convective

Q: what formula would give the vertical flux density of water vapour  $[\text{kg m}^{-2} \text{s}^{-1}]$ ?

"absolute humidity"

$w \rho_v$

- conduction/diffusion mechanism for transport is of negligible importance in the interior of the atmosphere
- but may dominate within thin layers (mm to cm in depth) adjacent to surfaces



## Lecture of 26 Sept.

- temperature change versus pressure change (Poisson's equation)
- relative importance of conductive/convective/radiative fluxes in soil/atmosphere/space
- laws for conductive and convective heat fluxes (i.e. eqns. quantifying such fluxes)
- how an appropriate conservation law for heat energy is "built" in the case of a homogeneous soil (a more complex law would be needed if we had to handle exchange of heat along the horizontal axes)
- role of symmetry in science laws – here an *assumed* symmetry allowed us to remove two of the four space-time coordinates, i.e. we assumed  $T$  is "homogeneous" along  $x,y$