

EAS 471

MODERN LS MODEL(S)

- ① Zeroth order 1D LSM is the RDM — Z is Markovian

$$dZ = (\partial K / \partial Z) dt + \sqrt{2K dt} \, r, \quad r \in N(0, 1)$$

Equivalent to eddy diffusion approach, therefore not rigorous in near-field of source(s).

- ② First order LSM is the generalized Langevin model, $dW = a dt + b dz$ — (Z, W) is a (compound) Markovian variable

- ③ Second order LSM is a "random walk in accel", $dA = a dt + b dz$ — (A, W, Z) jointly Markovian

We'll focus on ②, option for assignment 3.

Let $dU_i = (dU, dV, dW)$ be the increments in Lagr. velocity (particle velocity) over time step dt , where dt is "small" w.r.t. all signif. time scales

$$\begin{aligned} dZ &= W dt \\ dW &= a dt + b d\xi \end{aligned}$$

most general form of Markovian model for (W, Z)
 $d\xi \in N(0, dt)$

1D case

(i) To determine the coefft "b", invoke Kolmogorov's postulate that

$$\langle dU_i dU_j \rangle = C_0 \epsilon dt \delta_{ij}, \quad C_0 \text{ dimensionless const.}$$

ϵ turb. kinetic energy dissipation rate

Implies $\langle dW^2 \rangle = C_0 \epsilon dt$

$$\begin{aligned} &= \langle a^2 dt^2 + b^2 d\xi^2 + 2ab dt d\xi \rangle \\ &= \underbrace{\langle a^2 \rangle dt^2}_{\text{neglect}} + b^2 \underbrace{\langle d\xi^2 \rangle}_{dt} + 2b dt \underbrace{\langle a d\xi \rangle}_{\text{zero as } d\xi \text{ is chosen randomly}} \end{aligned}$$

$\therefore C_0 \epsilon dt = b^2 dt \quad \text{or} \quad b = \sqrt{C_0 \epsilon}$

(ii) How to get the other coefficient? Thomson's (1987) "well mixed condition"

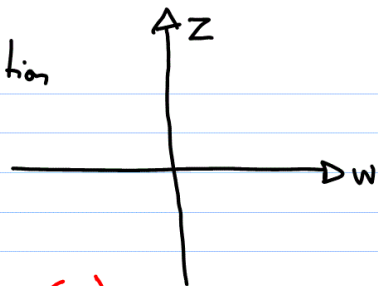
(3)

Let $p(z, w, t)$ be the joint PDF for ^{particle} vert. veloc. + position

Imagine that at $t=0$ we had particles

"well mixed" in the flow. (I'll restrict to constant

density flow). Let the Eulerian velocity pdf be $g_a(w)$.



$p(z, w, 0) \propto \rho g_a(w)$, this is the maximally disordered (max entropy) state

Therefore $p(z, w, dt)$ must also be well mixed. But p is known to evolve according to the Fokker-Plank eqn., which for this case reads

$$\begin{aligned} \frac{\partial p}{\partial t} &= - \frac{\partial}{\partial z} (w p) - \frac{\partial}{\partial w} (a p) + \frac{b^2}{2} \frac{\partial^2 p}{\partial w^2} \\ &= - \frac{\partial}{\partial z} (w p) - \frac{\partial}{\partial w} \left(a p - \frac{b^2}{2} \frac{\partial p}{\partial w} \right) \end{aligned}$$

(4)

Thomson's w.m.c. implies that $g_a(w)$ must be a steady solution of this F.P. equation. This gives a constraint on the coefficient a .

Vertically inhomogeneous Gaussian turbulence

$$g_a(w) = \frac{1}{\sqrt{2\pi} \sigma_w(z)} \exp\left[-\frac{w^2}{2\sigma_w^2}\right] \quad \text{Eulerian velocity PDF}$$

(we'll assume horiz. homog. flow with $\nabla \cdot \vec{u} = 0$)

$$\text{i.e. } \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0 \Rightarrow \frac{\cancel{\partial \bar{u}}}{\cancel{\partial x}} + \frac{\cancel{\partial \bar{v}}}{\cancel{\partial y}} + \frac{\partial \bar{w}}{\partial z} = 0$$

$\downarrow$ $\downarrow$
 hh hh

$$\bar{w}(z) = \bar{w}(0) = 0$$

$$a = -\frac{W}{\tau} + \frac{1}{2} \frac{\partial \sigma_w^2}{\partial z} \left(\frac{W^2}{\sigma_w^2} + 1 \right), \quad \tau = \frac{2\sigma_w^2}{C_0 \varepsilon}$$

⑤

hh ASL

$$\begin{aligned} \bar{u}(z) \\ \sigma_w(z) \\ \varepsilon(z) \end{aligned}$$

$$dX = \bar{u}(Z) dt$$

$$dt = \mu \tau = \frac{2}{c_0} \frac{\sigma_w^2(Z)}{\varepsilon(Z)}$$

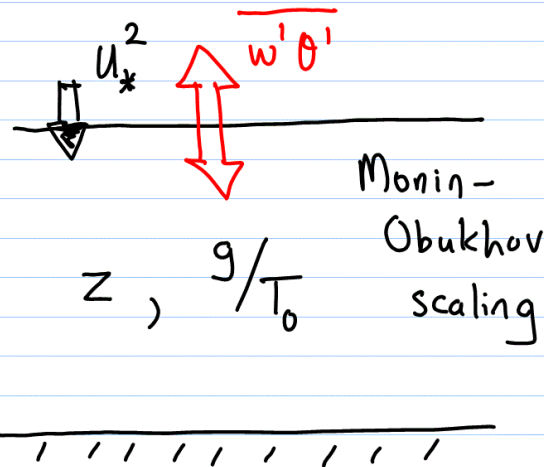
$$dZ = W dt$$

$$dW = a(Z, W) dt + \sqrt{C_0 \varepsilon(Z)} d\xi$$

ABOUT THE IDEAL HORIZ. HOMOG ASL *

Do a dimensional analysis for the property $\partial \bar{u} / \partial z$ ($n=5, m=3$)

* shallow layer $Z_0 \ll z \ll \delta$
 sfc roughness ABL depth



(6)

across which height variation of the (kinematic) fluxes of sensible heat and horiz. momentum can be neglected.

Monin & Obukhov conjectured that properties in the ASL "scale with" $z, g/T_0, u_*, \overline{w'\theta'}$ where T_0 is a bulk Kelvin temp.

$$\frac{\partial \bar{u}}{\partial z} = \sum_i c_i z^{\alpha_i} u_*^{\beta_i} (g/T_0)^{\gamma_i} (\overline{w'\theta'})^{\varepsilon_i} \quad \left. \vphantom{\frac{\partial \bar{u}}{\partial z}} \right\} \text{method of indices}$$



$$\frac{\partial \bar{u}}{\partial z} = \sum_i c_i z^{\gamma_i - 1} u_*^{-3\gamma_i + 1} (g/T_0)^{\gamma_i} \overline{w'\theta'}^{\gamma_i}$$

$$\Rightarrow \frac{z}{u_*} \frac{\partial \bar{u}}{\partial z} = \sum_i c_i \left(\frac{z}{\hat{L}} \right)^{\gamma_i} = f\left(\frac{z}{\hat{L}}\right), \quad \hat{L} = \frac{u_*^3}{(g/T_0) \overline{w'\theta'}}$$